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# Reconstruction of a potential from the impedance boundary map

M.I. Isaev and R.G. Novikov

## Abstract

We consider the inverse boundary value problem for the Schrödinger equation at fixed energy with boundary measurements represented as the impedance boundary map (or Robin-to-Robin map). We give formulas and equations for finding (generalized) scattering data for the aforementioned equation from boundary measurements in this impedance representation. Combining these results with results of the inverse scattering theory we obtain efficient methods for reconstructing potential from the impedance boundary map. To our knowledge, results of the present work are new already for the case of Neumann-to-Dirichlet map.

## 1 Introduction

We consider the equation

$$-\Delta\psi + v(x)\psi = E\psi, \quad x \in D, \quad E \in \mathbb{R}, \quad (1.1)$$

where

$$\begin{aligned} D \text{ is an open bounded domain in } \mathbb{R}^d, \quad d \geq 2, \\ \text{with } \partial D \in C^2, \end{aligned} \quad (1.2)$$

$$v \in \mathbb{L}^\infty(D), \quad v = \bar{v}. \quad (1.3)$$

Equation (1.1) can be considered as the stationary Schrödinger equation of quantum mechanics at fixed energy  $E$ . Equation (1.1) at fixed  $E$  arises also in acoustics and electrodynamics.

Following [19], [26], we consider the impedance boundary map  $\hat{M}_\alpha = \hat{M}_{\alpha,v}(E)$  defined by

$$\hat{M}_\alpha[\psi]_\alpha = [\psi]_{\alpha-\pi/2} \quad (1.4)$$

for all sufficiently regular solutions  $\psi$  of equation (1.1) in  $\bar{D} = D \cup \partial D$ , where

$$[\psi]_\alpha = [\psi(x)]_\alpha = \cos \alpha \psi(x) - \sin \alpha \frac{\partial \psi}{\partial \nu}|_{\partial D}(x), \quad x \in \partial D, \quad \alpha \in \mathbb{R} \quad (1.5)$$

and  $\nu$  is the outward normal to  $\partial D$ . Under assumptions (1.2), (1.3), in Lemma 3.2 of [26] it was shown that there is not more than a countable number of  $\alpha \in \mathbb{R}$  such that  $E$  is an eigenvalue for the operator  $-\Delta + v$  in  $D$  with the boundary condition

$$\cos \alpha \psi|_{\partial D} - \sin \alpha \frac{\partial \psi}{\partial \nu}|_{\partial D} = 0. \quad (1.6)$$

Therefore, for any fixed  $E$  we can assume that for some fixed  $\alpha \in \mathbb{R}$

$$\begin{aligned} E \text{ is not an eigenvalue for the operator } -\Delta + v \text{ in } D \\ \text{with boundary condition (1.6)} \end{aligned} \quad (1.7)$$

and, as a corollary,  $\hat{M}_\alpha$  can be defined correctly.

We consider  $\hat{M}_\alpha = \hat{M}_{\alpha,v}(E)$  as an operator representation of all possible boundary measurements for the physical model described by (1.1). We recall that the impedance boundary map  $\hat{M}_\alpha$  is reduced to the Dirichlet-to-Neumann(DtN) map if  $\alpha = 0$  and is reduced to the Neumann-to-Dirichlet(NtD) map if  $\alpha = \pi/2$ . The map  $\hat{M}_\alpha$  can be called also as the Robin-to-Robin map.

As in [26], we consider the following inverse boundary value problem for equation (1.1).

**Problem 1.1.** Given  $\hat{M}_\alpha$  for some fixed  $E$  and  $\alpha$ , find  $v$ .

This problem can be considered as the Gel'fand inverse boundary value problem for the Schrödinger equation at fixed energy (see [18], [36]). Note that in the initial Gel'fand formulation energy  $E$  was not yet fixed and boundary measurements were considered as an operator relating  $\psi|_{\partial D}$  and  $\frac{\partial \psi}{\partial \nu}|_{\partial D}$  for  $\psi$  satisfying (1.1).

Problem 1.1 for  $E = 0$  can be considered also as a generalization of the Calderon problem of the electrical impedance tomography (see [13], [36]).

Note also that Problem 1.1 can be considered as an example of ill-posed problem: see [4], [31] for an introduction to this theory.

Problem 1.1 includes, in particular, the following questions: (a) uniqueness, (b) reconstruction, (c) stability.

Global uniqueness theorems and global reconstruction methods for Problem 1.1 with  $\alpha = 0$  (i.e. for the DtN case) were given for the first time in [36] in dimension  $d \geq 3$  and in [10] in dimension  $d = 2$ .

Global stability estimates for Problem 1.1 with  $\alpha = 0$  were given for the first time in [1] in dimension  $d \geq 3$  and in [45] in dimension  $d = 2$ . A principal improvement of the result of [1] was given recently in [44] (for  $E = 0$ ). Due to [32] these logarithmic stability results are optimal (up to the value of the exponent). An extension of the instability estimates of [32] to the case of non-zero energy as well as to the case of Dirichlet-to-Neumann map given on the energy intervals was obtained in [24]. An extension of stability estimates of [44] to the energy dependent case was given recently in [27]. Instability estimates complementing stability results of [27] were obtained in [25].

Note also that for the Calderon problem (of the electrical impedance tomography) in its initial formulation the global uniqueness was firstly proved in [51] for  $d \geq 3$  and in [35] for  $d = 2$ . In addition, for the case of piecewise constant or piecewise real analytic conductivity the first uniqueness results for the Calderon problem in dimension  $d \geq 2$  were given in [15], [28]. Lipschitz stability estimate for the case of piecewise constant conductivity was proved in [2] and additional studies in this direction were fulfilled in [48].

It should be noted that in most of previous works on inverse boundary value problems for equation (1.1) at fixed  $E$  it was assumed in one way or another that  $E$  is not a Dirichlet eigenvalue for the operator  $-\Delta + v$  in  $D$ , see [1], [32], [36], [44]-[49]. Nevertheless, the results of [10] can be considered as global uniqueness and reconstruction results for Problem 1.1 in dimension  $d = 2$  with general  $\alpha$ .

Global stability estimates for Problem 1.1 in dimension  $d \geq 2$  with general  $\alpha$  were recently given in [26].

In the present work we give formulas and equations for finding (generalized) scattering data from the impedance boundary map  $\hat{M}_\alpha$  with general  $\alpha$ . Combining these results with results of [21], [23], [35], [37]-[39], [41]-[43], we obtain efficient reconstruction methods for Problem 1.1 in multidimensions with general  $\alpha$ . To our knowledge these results are new already for the NtD case.

In particular, in the present work we give first mathematically justified approach for reconstructing coefficient  $v$  from boundary measurements for (1.1) via inverse scattering without the assumption that  $E$  is not a Dirichlet

eigenvalue for  $-\Delta + v$  in  $D$ . In addition, numerical efficiency of related inverse scattering techniques was shown in [3], [9], [11], [12]; see also [8].

Definitions of (generalized) scattering data are recalled in Section 2. Our main results are presented in Section 3. Proofs of these results are given in Sections 4, 5 and 6.

## 2 Scattering data

Consider the Schrödinger equation

$$-\Delta\psi + v(x)\psi = E\psi, \quad x \in \mathbb{R}^d, \quad d \geq 2 \quad (2.1)$$

where

$$(1 + |x|)^{d+\varepsilon}v(x) \in \mathbb{L}^\infty(\mathbb{R}^d) \quad (\text{as a function of } x), \quad \text{for some } \varepsilon > 0. \quad (2.2)$$

For equation (2.1) we consider the functions  $\psi^+$  and  $f$  of the classical scattering theory and the Faddeev functions  $\psi$ ,  $h$ ,  $\psi_\gamma$ ,  $h_\gamma$  (see, for example, [6], [14], [16], [17], [20], [23], [33], [37]).

The functions  $\psi^+$  and  $f$  can be defined as follows:

$$\psi^+(x, k) = e^{ikx} + \int_{\mathbb{R}^d} G^+(x - y, k) v(y) \psi^+(y, k) dy, \quad (2.3)$$

$$G^+(x, k) = - \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} \frac{e^{i\xi x}}{\xi^2 - k^2 - i0} d\xi, \quad (2.4)$$

$$x, k \in \mathbb{R}^d, \quad k^2 > 0,$$

where (2.3) at fixed  $k$  is considered as an equation for  $\psi^+$  in  $\mathbb{L}^\infty(\mathbb{R}^d)$ ;

$$f(k, l) = \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} e^{-ilx} \psi^+(x, k) v(x) dx, \quad (2.5)$$

$$k, l \in \mathbb{R}^d, \quad k^2 > 0.$$

In addition:  $\psi^+(x, k)$  satisfies (2.3) for  $E = k^2$  and describes scattering of the plane waves  $e^{ikx}$ ;  $f(k, l)$ ,  $k^2 = l^2$ , is the scattering amplitude for equation (2.1) for  $E = k^2$ . Equation (2.3) is the Lippman-Schwinger integral equation.

The functions  $\psi$  and  $h$  can be defined as follows:

$$\psi(x, k) = e^{ikx} + \int_{\mathbb{R}^d} G(x - y, k) v(y) \psi(y, k) dy, \quad (2.6)$$

$$G(x, k) = - \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} \frac{e^{i\xi x} d\xi}{\xi^2 + 2k\xi} e^{ikx}, \quad (2.7)$$

$$x \in \mathbb{R}^d, \quad k \in \mathbb{C}^d, \quad \text{Im } k \neq 0,$$

where (2.6) at fixed  $k$  is considered as an equation for  $\psi = e^{ikx} \mu(x, k)$ ,  $\mu \in \mathbb{L}^\infty(\mathbb{R}^d)$ ;

$$h(k, l) = \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} e^{-ilx} \psi(x, k) v(x) dx, \quad (2.8)$$

$$k, l \in \mathbb{C}^d, \quad \text{Im } k = \text{Im } l \neq 0.$$

In addition,  $\psi(x, k)$  satisfies (2.1) for  $E = k^2$ , and  $\psi$ ,  $G$  and  $h$  are (nonanalytic) continuations of  $\psi^+$ ,  $G^+$  and  $f$  to the complex domain. In particular,  $h(k, l)$  for  $k^2 = l^2$  can be considered as the "scattering" amplitude in the complex domain for equation (2.1) for  $E = k^2$ . The functions  $\psi_\gamma$  and  $h_\gamma$  are defined as follows:

$$\begin{aligned} \psi_\gamma(x, k) &= \psi(x, k + i0\gamma), & h_\gamma(k, l) &= h(k + i0\gamma, l + i0\gamma), \\ x, k, l, \gamma &\in \mathbb{R}^d, \quad \gamma^2 = 1. \end{aligned} \quad (2.9)$$

We recall also that

$$\begin{aligned} \psi^+(x, k) &= \psi_{k/|k|}(x, k), & f(k, l) &= h_{k/|k|}(k, l), \\ x, k, l &\in \mathbb{R}^d, \quad |k| > 0. \end{aligned} \quad (2.10)$$

We consider  $f(k, l)$  and  $h_\gamma(k, l)$ , where  $k, l, \gamma \in \mathbb{R}^d$ ,  $k^2 = l^2 = E$ ,  $\gamma^2 = 1$ , and  $h(k, l)$ , where  $k, l \in \mathbb{C}^d$ ,  $\text{Im } k = \text{Im } l \neq 0$ ,  $k^2 = l^2 = E$ , as scattering data  $S_E$  for equation (2.1) at fixed  $E \in (0, +\infty)$ . We consider  $h(k, l)$ , where  $k, l \in \mathbb{C}^d$ ,  $\text{Im } k = \text{Im } l \neq 0$ ,  $k^2 = l^2 = E$ , as scattering data  $S_E$  for equation (2.1) at fixed  $E \in (-\infty, 0]$ .

We consider also the sets  $\mathcal{E}$ ,  $\mathcal{E}_\gamma$ ,  $\mathcal{E}^+$  defined as follows:

$$\mathcal{E} = \left\{ \zeta \in \mathbb{C}^d \setminus \mathbb{R}^d : \begin{array}{l} \text{equation (2.6) for } k = \zeta \text{ is not} \\ \text{uniquely solvable for } \psi = e^{ikx} \mu \text{ with } \mu \in \mathbb{L}^\infty(\mathbb{R}^d) \end{array} \right\}, \quad (2.11a)$$

$$\mathcal{E}_\gamma = \left\{ \begin{array}{l} \zeta \in \mathbb{R}^d \setminus \{0\} : \text{equation (2.6) for } k = \zeta + i0\gamma \\ \text{is not uniquely solvable for } \psi = \mathbb{L}^\infty(\mathbb{R}^d) \end{array} \right\}, \quad (2.11b)$$

$$\gamma \in \mathbb{S}^{d-1},$$

$$\mathcal{E}^+ = \left\{ \begin{array}{l} \zeta \in \mathbb{R}^d \setminus \{0\} : \text{equation (2.6) for } k = \zeta \text{ is not} \\ \text{uniquely solvable for } \psi = \mathbb{L}^\infty(\mathbb{R}^d) \end{array} \right\}. \quad (2.11c)$$

In addition,  $\mathcal{E}^+$  is a well-known set of the classical scattering theory for equation (2.1) and  $\mathcal{E}^+ = \emptyset$  for real-valued  $v$  satisfying (2.2) (see, for example, [6], [33]). Note also that  $\mathcal{E}^+$  is spherically symmetric. The sets  $\mathcal{E}$ ,  $\mathcal{E}_\gamma$  were considered for the first time in [16], [17]. Concerning the properties of  $\mathcal{E}$  and  $\mathcal{E}_\gamma$ , see [17], [22], [23], [30], [33], [35], [38], [52].

We consider also the functions  $R$ ,  $R_\gamma$ ,  $R^+$  defined as follows:

$$R(x, y, k) = G(x - y, k) + \int_{\mathbb{R}^d} G(x - z, k) v(z) R(z, y, k) dz, \quad (2.12)$$

$$x, y \in \mathbb{R}^d, \quad k \in \mathbb{C}^d, \quad \text{Im } k \neq 0,$$

where  $G$  is defined by (2.7) and formula (2.12) at fixed  $y$ ,  $k$  is considered as an equation for

$$R(x, y, k) = e^{ik(x-y)} r(x, y, k), \quad (2.13)$$

where  $r$  is sought with the properties

$$r(\cdot, y, k) \text{ is continuous on } \mathbb{R}^d \setminus \{y\} \quad (2.14a)$$

$$r(x, y, k) \rightarrow 0 \text{ as } |x| \rightarrow \infty, \quad (2.14b)$$

$$\begin{aligned} r(x, y, k) &= O(|x - y|^{2-d}) \text{ as } x \rightarrow y \text{ for } d \geq 3, \\ r(x, y, k) &= O(|\ln |x - y||) \text{ as } x \rightarrow y \text{ for } d = 2; \end{aligned} \quad (2.14c)$$

$$\begin{aligned} R_\gamma(x, y, k) &= R(x, y, k + i0\gamma), \\ x, y &\in \mathbb{R}^d, \quad k \in \mathbb{R}^d \setminus \{0\}, \quad \gamma \in \mathbb{S}^{d-1}; \end{aligned} \quad (2.15)$$

$$\begin{aligned} R^+(x, y, k) &= R_{k/|k|}(x, y, k), \\ x, y &\in \mathbb{R}^d, \quad k \in \mathbb{R}^d \setminus \{0\}. \end{aligned} \quad (2.16)$$

In addition, the functions  $R(x, y, k)$ ,  $R_\gamma(x, y, k)$  and  $R^+(x, y, k)$  (for their domains of definition in  $k$  and  $\gamma$ ) satisfy the following equations:

$$\begin{aligned}(\Delta_x + E - v(x))R(x, y, k) &= \delta(x - y), \\(\Delta_y + E - v(y))R(x, y, k) &= \delta(x - y), \\x, y &\in \mathbb{R}^d, \quad E = k^2.\end{aligned}\tag{2.17}$$

The function  $R^+(x, y, k)$  (defined by means of (2.12) for  $k \in \mathbb{R}^d \setminus \{0\}$  with  $G$  replaced by  $G^+$  of (2.4)) is well-known in the scattering theory for equations (2.1), (2.17) (see, for example, [7]). In particular, this function describes scattering of the spherical waves  $G^+(x - y, k)$  generated by a source at  $y$ . In addition  $R^+(x, y, k)$  is a radial function in  $k$ , i.e.

$$R^+(x, y, k) = R^+(x, y, |k|), \quad x, y \in \mathbb{R}^d, \quad k \in \mathbb{R}^d \setminus \{0\}.\tag{2.18}$$

Apparently, the functions  $R$  and  $R_\gamma$  were considered for the first time in [38].

In addition, under the assumption (2.2): equation (2.12) at fixed  $y$  and  $k$  is uniquely solvable for  $R$  with the properties (2.13), (2.14) if and only if  $k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E})$ ; equation (2.12) with  $k = \zeta + i0\gamma$ ,  $\zeta \in \mathbb{R}^d \setminus \{0\}$ ,  $\gamma \in \mathbb{S}^{d-1}$ , at fixed  $y$ ,  $\zeta$  and  $\gamma$  is uniquely solvable for  $R_\gamma$  if and only if  $\zeta \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma)$ ; equation (2.12) with  $k = \zeta + i0\zeta/|\zeta|$ ,  $\zeta \in \mathbb{R}^d \setminus 0$ , at fixed  $y$  and  $\zeta$  is uniquely solvable for  $R^+$  if and only if  $\zeta \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^+)$ .

### 3 Main results

Let  $v$  and  $v^0$  satisfy (1.3), (1.7) for some fixed  $E$  and  $\alpha$ . Let  $M_{\alpha,v}(x, y, E)$ ,  $M_{\alpha,v^0}(x, y, E)$ ,  $x, y \in \partial D$ , denote the Schwartz kernels of the impedance boundary maps  $\hat{M}_{\alpha,v}$ ,  $\hat{M}_{\alpha,v^0}$ , for potentials  $v$  and  $v^0$ , respectively, where  $\hat{M}_{\alpha,v}$ ,  $\hat{M}_{\alpha,v^0}$  are considered as linear integral operators. In addition, we consider  $v^0$  as some known background potential.

Let  $h, \psi, f, \psi^+, h_\gamma, \psi_\gamma, \mathcal{E}, \mathcal{E}^+, \mathcal{E}_\gamma$  and  $h^0, \psi^0, f^0, \psi^{+,0}, h_\gamma^0, \psi_\gamma^0, \mathcal{E}^0, \mathcal{E}^{+,0}, \mathcal{E}_\gamma^0$  denote the functions and sets of (2.3), (2.5), (2.6), (2.8), (2.9), (2.11) for potentials  $v$  and  $v^0$ , respectively. Here and below in this section we always assume that  $v \equiv 0$ ,  $v^0 \equiv 0$  on  $\mathbb{R}^d \setminus D$ .



**Theorem 3.1.** *Let  $D$  satisfy (1.2) and potentials  $v, v^0$  satisfy (1.3), (1.7) for some fixed  $E$  and  $\alpha$ . Then:*

$$\begin{aligned} h(k, l) - h^0(k, l) &= \\ &= \left(\frac{1}{2\pi}\right)^d \int_{\partial D} \int_{\partial D} [\psi^0(x, -l)]_\alpha (M_{\alpha, v} - M_{\alpha, v^0})(x, y, E) [\psi(y, k)]_\alpha dx dy, \\ &\quad k, l \in \mathbb{C}^d \setminus (\mathcal{E} \cup \mathcal{E}^0), \quad k^2 = l^2 = E, \quad \text{Im } k = \text{Im } l \neq 0, \end{aligned} \quad (3.1)$$

$$\begin{aligned} [\psi(x, k)]_\alpha &= [\psi^0(x, k)]_\alpha + \int_{\partial D} A_\alpha(x, y, k) [\psi(y, k)]_\alpha dy, \\ x \in \partial D, \quad k \in \mathbb{C}^d \setminus (\mathcal{E} \cup \mathcal{E}^0), \quad \text{Im } k \neq 0, \quad k^2 = E \end{aligned} \quad (3.2)$$

where

$$A_\alpha(x, y, k) = \lim_{\varepsilon \rightarrow +0} \int_{\partial D} D_{\alpha, \varepsilon} R^0(x, \xi, k) (M_{\alpha, v} - M_{\alpha, v^0})(\xi, y, E) d\xi, \quad (3.3)$$

$$\begin{aligned} D_{\alpha, \varepsilon} R^0(x, \xi, k) &= [[R^0(x + \varepsilon \nu_x, \xi, k)]_{\xi, \alpha}]_{x, \alpha} = \\ &= \left( \cos^2 \alpha - \sin \alpha \cos \alpha \left( \frac{\partial}{\partial \nu_x} + \frac{\partial}{\partial \nu_\xi} \right) + \sin^2 \alpha \frac{\partial^2}{\partial \nu_x \partial \nu_\xi} \right) R^0(x + \varepsilon \nu_x, \xi, k), \\ &\quad x, \xi, y \in \partial D, \end{aligned} \quad (3.4)$$

where  $R^0$  denotes the Green function of (2.12) for potential  $v^0$ ,  $\nu_x$  is the outward normal to  $\partial D$  at  $x$ . In addition, formulas completely similar to (3.1) - (3.4) are also valid for the classical scattering functions  $f, \psi^+, f^0, \psi^{+,0}$  and sets  $\mathcal{E}^+, \mathcal{E}^{+,0}$  of (2.3), (2.5), (2.11c) for  $v$  and  $v^0$ , respectively, but with  $R^{+,0}$  in place of  $R^0$  in (3.3), (3.4), where  $R^{+,0}$  denotes the Green function of (2.16) for potential  $v^0$ .

Theorem 3.1 is proved in Section 4.

Note that formula of the type (3.1) for  $h_\gamma$  is not completely similar to (3.1): see formula (3.6) given below. In this formula (3.6), in addition to expected  $\psi_\gamma(x, k)$ , we use also  $\psi_\gamma(x, k, l)$  defined as follows:

$$\begin{aligned}
\psi_\gamma(x, k, l) &= e^{ilx} + \int_{\mathbb{R}^d} G_\gamma(x - y, k) v(y) \psi_\gamma(y, k, l) dy, \\
G_\gamma(x, k) &= G(x, k + i0\gamma), \\
\gamma &\in \mathbb{S}^{d-1}, \quad x, k, l \in \mathbb{R}^d, \quad k^2 = l^2 > 0,
\end{aligned} \tag{3.5}$$

where (3.5) at fixed  $\gamma, k, l$  is considered as an equation for  $\psi_\gamma(\cdot, k, l)$  in  $\mathbb{L}^\infty(\mathbb{R}^d)$ ,  $G$  is defined by (2.7).

**Proposition 3.1.** *Let the assumptions of Theorem 3.1 hold. Let  $\psi_\gamma(x, k)$  correspond to  $v$  according to (2.9) and  $\psi_{-\gamma}^0(\cdot, k, l)$  correspond to  $v^0$  according to (3.5). Then*

$$\begin{aligned}
h_\gamma(k, l) - h_\gamma^0(k, l) &= \\
&= \left( \frac{1}{2\pi} \right)^d \int_{\partial D} \int_{\partial D} [\psi_{-\gamma}^0(x, -k, -l)]_\alpha (M_{\alpha, v} - M_{\alpha, v^0})(x, y, E) [\psi_\gamma(y, k)]_\alpha dx dy, \\
&\quad \gamma \in \mathbb{S}^{d-1}, \quad k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma \cup \mathcal{E}_\gamma^0), \quad l \in \mathbb{R}^d, \quad k^2 = l^2 = E.
\end{aligned} \tag{3.6}$$

In addition, formulas completely similar to (3.2) - (3.4) are also valid for the functions  $\psi_\gamma(x, k)$ ,  $\psi_\gamma^0(x, k)$  and sets  $\mathcal{E}_\gamma$ ,  $\mathcal{E}_\gamma^0$  of (2.9), (2.11b) for  $v$  and  $v^0$ , respectively, but with  $R_\gamma^0$  in place of  $R^0$  in (3.3), (3.4), where  $R_\gamma^0$  denotes the Green function of (2.15) for potential  $v^0$ .

Proposition 3.1 is proved in Section 4.

Note that (3.2) is considered as a linear integral equation for finding  $[\psi(x, k)]_\alpha$ ,  $x \in \partial D$ , at fixed  $k$ , from  $\hat{M}_{\alpha, v} - \hat{M}_{\alpha, v^0}$  and  $[\psi^0(x, k)]_\alpha$ , whereas (3.1) is considered as an explicit formula for finding  $h$  from  $h^0$ ,  $\hat{M}_{\alpha, v} - \hat{M}_{\alpha, v^0}$ ,  $[\psi^0(x, k)]_\alpha$  and  $[\psi(x, k)]_\alpha$ . In addition, we use similar interpretation for similar formulas for  $\psi^+$ ,  $f$  and for  $\psi_\gamma$ ,  $h_\gamma$ , mentioned in Theorem 3.1 and Proposition 3.1.

Under the assumptions of Theorem 3.1, the following propositions are valid:

**Proposition 3.2.** *Equation (3.2) for  $[\psi(x, k)]_\alpha$  at fixed  $k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}^0)$  is a Fredholm linear integral equation of the second kind in the space of bounded functions on  $\partial D$ . In addition, the same is also valid for the equation for  $[\psi^+(x, k)]_\alpha$  at fixed  $k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^{+, 0})$ , mentioned in Theorem 3.1, and for the equation for  $[\psi_\gamma(x, k)]_\alpha$  at fixed  $\gamma \in \mathbb{S}^{d-1}$ ,  $k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma^0)$ , mentioned in Proposition 3.1.*

Proposition 3.2 is proved in Section 4.

**Proposition 3.3.** *For  $k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}^0)$  equation (3.2) is uniquely solvable in the space of bounded functions on  $\partial D$  if and only if  $k \notin \mathcal{E}$ . In addition, the aforementioned equations for  $[\psi^+(x, k)]_\alpha$ ,  $k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^{+,0})$ , and  $[\psi_\gamma(x, k)]_\alpha$ ,  $\gamma \in \mathbb{S}^{d-1}$ ,  $k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma^0)$ , are uniquely solvable in the space of bounded functions on  $\partial D$  if and only if  $k \notin \mathcal{E}^+$  and  $k \notin \mathcal{E}_\gamma$ , respectively.*

Proposition 3.3 is proved in Section 5.

**Proposition 3.4.** *Let  $\phi_\alpha(x, y)$  be the solution of the Dirichlet boundary value problem at fixed  $y \in \partial D$ ,  $\lambda \in \mathbb{C}$ :*

$$\begin{aligned} -\Delta_x \phi_\alpha(x, y) &= \lambda \phi_\alpha(x, y), & x \in D, \\ \phi_\alpha(x, y) &= (M_{\alpha, v} - M_{\alpha, v^0})(x, y, E), & x \in \partial D, \end{aligned} \quad (3.7)$$

where we assume that  $\lambda$  is not a Dirichlet eigenvalue for  $-\Delta$  in  $D$ . Then

$$\begin{aligned} A_\alpha(x, y, k) &= \lim_{\varepsilon \rightarrow +0} \int_{\partial D} [R^0(x + \varepsilon \nu_x, \xi, k)]_{x, \alpha} [\phi_\alpha(\xi, y)]_{\xi, \alpha} d\xi - \\ &\quad - \sin \alpha \int_D [R^0(x, \xi, k)]_{x, \alpha} (v^0(\xi) - E + \lambda) \phi_\alpha(\xi, y) d\xi, \quad x, y \in \partial D, \end{aligned} \quad (3.8)$$

where

$$\begin{aligned} [R^0(x + \varepsilon \nu_x, \xi, k)]_{x, \alpha} &= \left( \cos \alpha - \sin \alpha \frac{\partial}{\partial \nu_x} \right) R^0(x + \varepsilon \nu_x, \xi, k), \quad x \in \partial D, \xi \in \bar{D}, \\ [\phi_\alpha(\xi, y)]_{\xi, \alpha} &= \left( \cos \alpha - \sin \alpha \frac{\partial}{\partial \nu_\xi} \right) \phi_\alpha(\xi, y) = \\ &= \cos \alpha \phi_\alpha(\xi, y) - \sin \alpha \left( \hat{\Phi}(\lambda) \phi_\alpha(\cdot, y) \right)(\xi), \quad \xi, y \in \partial D, \end{aligned} \quad (3.9)$$

where  $A_\alpha$  is defined in (3.3),  $\hat{\Phi}(\lambda) = \hat{M}_{0,0}(\lambda)$  is the Dirichlet-to-Neumann map for (3.7). In addition, formulas completely similar to (3.8) are also valid for the kernels  $A_\alpha^+$  (but with  $R_0^+$  in place of  $R_0$ ) and  $A_{\alpha, \gamma}$  (but with  $R_\gamma^0$  in place of  $R^0$ ), arising in the equations for  $[\psi^+]_\alpha$  and  $[\psi_\gamma]_\alpha$ , mentioned in Theorem 3.1 and Proposition 3.1.

Proposition 3.4 is proved in Section 4.

Note that, for the case when  $\sin \alpha = 0$ , formula (3.8) coincides with (3.3). However, for  $\sin \alpha \neq 0$ , formula (3.8) does not contain  $\partial^2 R^0 / \partial \nu_x \partial \nu_\xi$  in contrast with (3.3) and is more convenient than (3.8) in this sense.

Theorem 3.1, Propositions 3.1 - 3.4 and the reconstruction results from generalized scattering data (see [20], [21], [23], [37]- [39], [41]- [43], [47]) imply the following corollary:

**Corollary 3.1.** To reconstruct a potential  $v$  in the domain  $D$  from its impedance boundary map  $\hat{M}_{\alpha,v}(E)$  at fixed  $E$  and  $\alpha$  one can use the following schema:

1.  $v^0 \rightarrow \{S_E^0\}, \{R^0\}, \{[\psi^0]_\alpha\}, \hat{M}_{\alpha,v^0}$  via direct problem methods,
2.  $\{R^0\}, \hat{M}_{\alpha,v^0}, \hat{M}_{\alpha,v} \rightarrow \{A_\alpha\}$  as described in Theorem 3.1 and Propositions 3.1, 3.4,
3.  $\{A_\alpha\}, \{[\psi^0]_\alpha\} \rightarrow \{[\psi]_\alpha\}$  as described in Theorem 3.1 and Proposition 3.1,
4.  $\{S_E^0\}, \{[\psi^0]_\alpha\}, \{[\psi]_\alpha\}, \hat{M}_{\alpha,v^0}, \hat{M}_{\alpha,v} \rightarrow \{S_E\}$  as described in Theorem 3.1 and Proposition 3.1,
5.  $\{S_E\} \rightarrow v$  as described in [20], [21], [23], [37]- [39], [41]- [43], [47],

where  $\{S_E^0\}$  and  $\{S_E\}$  denote some appropriate part of  $h^0, f^0, h_\gamma^0$  and  $h, f, h_\gamma$ , respectively,  $\{[\psi^0]_\alpha\}$  and  $\{[\psi]_\alpha\}$  denote some appropriate part of  $[\psi^0]_\alpha, [\psi^{+,0}]_\alpha, [\psi_\gamma^0]_\alpha$  and  $[\psi]_\alpha, [\psi^+]_\alpha, [\psi_\gamma]_\alpha$ , respectively,  $\{R^0\}, \{A_\alpha\}$  denote some appropriate part of  $R^0, R^{+,0}, R_\gamma^0, A_\alpha, A_\alpha^+, A_{\alpha,\gamma}$ .

**Remark 3.1.** For the case when  $v^0 \equiv 0, \sin \alpha = 0$ , Theorem 3.1, Propositions 3.1 - 3.3 and Corollary 3.1 (with available references at that time at step 5) were obtained in [36] (see also [34], [35]). Note that basic results of [36] were presented already in the survey given in [23]. For the case when  $\sin \alpha = 0$  Theorem 3.1, Propositions 3.1 - 3.3 and Corollary 3.1 (with available references at that time at step 5) were obtained in [40].

**Remark 3.2.** The results of Theorem 3.1, Propositions 3.1 - 3.4 and Corollary 3.1 remain valid for complex-valued  $v, v^0$  and complex  $E, \alpha$ , under the condition that (1.7) holds for both  $v$  and  $v^0$ .

**Remark 3.3.** Under the assumptions of Theorem 3.1, the following formula holds:

$$\hat{M}_{\alpha,v}(E) - \hat{M}_{\alpha,v^0}(E) = (D_\alpha R^{+,0}(E))^{-1} - (D_\alpha R^+(E))^{-1}, \quad (3.11)$$

$$\begin{aligned} D_\alpha R^+(E)u(x) &= \lim_{\varepsilon \rightarrow +0} \int_{\partial D} D_{\alpha,\varepsilon} R^+(x, y, \sqrt{E}) u(y) dy, \\ D_\alpha R^{+,0}(E)u(x) &= \lim_{\varepsilon \rightarrow +0} \int_{\partial D} D_{\alpha,\varepsilon} R^{+,0}(x, y, \sqrt{E}) u(y) dy, \end{aligned} \quad (3.12)$$

$$x \in \partial D,$$

where  $D_{\alpha,\varepsilon}$  is defined as in (3.4),  $R^+(x, y, \sqrt{E})$ ,  $R^{+,0}(x, y, \sqrt{E})$ ,  $\sqrt{E} > 0$ , are the Green functions of (2.16) written as in (2.18) for potentials  $v$ ,  $v^0$ , respectively,  $u$  is the test function. For the case when  $\sin \alpha = 0$ ,  $v^0 \equiv 0$ ,  $d \geq 3$ , formula (3.11) was given in [34]. Using techniques developed in [26] and in the present work, we obtain (3.11) in the general case.

## 4 Proofs of Theorem 3.1 and Propositions 3.1, 3.2, 3.4

In this section we will use formulas and equations for impedance boundary map from [26]. These results are presented in detail in Subsection 4.1. Proofs of Theorem 3.1 and Propositions 3.1, 3.2, 3.4 are given in Subsections 4.2, 4.3.

### 4.1 Preliminaries

Let  $G_{\alpha,v}(x, y, E)$  be the Green function for the operator  $\Delta - v + E$  in  $D$  with the impedance boundary condition (1.6) under assumptions (1.2), (1.3) and (1.7). We recall that (see formulas (3.12), (3.13) of [26]):

$$G_{\alpha,v}(x, y, E) = G_{\alpha,v}(y, x, E), \quad x, y \in \bar{D}, \quad (4.1)$$

and, for  $\sin \alpha \neq 0$ ,

$$M_{\alpha,v}(x, y, E) = \frac{1}{\sin^2 \alpha} G_{\alpha,v}(x, y, E) - \frac{\cos \alpha}{\sin \alpha} \delta_{\partial D}(x - y), \quad x, y \in \partial D, \quad (4.2)$$

where  $M_\alpha(x, y, E)$  and  $\delta_{\partial D}(x-y)$  denote the Schwartz kernels of the impedance boundary map  $\hat{M}_{\alpha,v}(E)$  and the identity operator  $\hat{I}$  on  $\partial D$ , respectively, where  $\hat{M}_\alpha$  and  $\hat{I}$  are considered as linear integral operators.

We recall also that (see, for example, formula (3.16) of [26]):

$$\psi(x) = \frac{1}{\sin \alpha} \int_{\partial D} (\cos \alpha \psi(\xi) - \sin \alpha \frac{\partial}{\partial \nu} \psi(\xi)) G_{\alpha,v}(x, \xi, E) d\xi, \quad x \in D, \quad (4.3)$$

for all sufficiently regular solutions  $\psi$  of equation (1.1) in  $\bar{D}$  and  $\sin \alpha \neq 0$ .

We will use the following properties of the Green function  $G_\alpha(x, y, E)$ :

$$G_{\alpha,v}(x, y, E) \text{ is continuous in } x, y \in \bar{D}, \quad x \neq y, \quad (4.4)$$

$$\begin{aligned} |G_{\alpha,v}(x, y, E)| &\leq c_1(|x - y|^{2-d}), \quad x, y \in \bar{D}, \text{ for } d \geq 3, \\ |G_{\alpha,v}(x, y, E)| &\leq c_1(|\ln |x - y||), \quad x, y \in \bar{D}, \text{ for } d = 2, \end{aligned} \quad (4.5)$$

where  $c_1 = c_1(D, E, v, \alpha) > 0$ .

Actually, properties (4.4), (4.5) are well-known for  $\sin \alpha = 0$  (the case of the Dirichlet boundary condition) and for  $\cos \alpha = 0$  (the case of the Neumann boundary condition). Properties (4.4), (4.5) with  $d \geq 3$ ,  $\frac{\sin \alpha}{\cos \alpha} < 0$ ,  $v \equiv 0$  and  $E = 0$  were proven in [29]. For  $d = 2$  see also [5]. In Section 6 we give proofs of (4.4), (4.5) for the case of general  $\alpha$ ,  $v$  and  $E$ .

In addition, under assumptions of Theorem 3.1, the following identity holds (see formula (3.9) of [26]):

$$\int_D (v - v^0) \psi \psi^0 dx = \int_{\partial D} [\psi]_\alpha \left( \hat{M}_{\alpha,v} - \hat{M}_{\alpha,v^0} \right) [\psi^0]_\alpha dx \quad (4.6)$$

for all sufficiently regular solutions  $\psi, \psi^0$  of equation (1.1) in  $\bar{D}$  for potentials  $v, v^0$ , respectively, where  $[\psi]_\alpha, [\psi^0]_\alpha$  are defined according to (1.5).

Identity (4.6) for  $\sin \alpha = 0$  is reduced to the Alessandrini identity (Lemma 1 of [1]).

We will use also that:

$$\begin{aligned} \|\hat{R}(k)u\|_{C^{1+\delta}(\Omega)} &\leq c_2(D, \Omega, v, k, \delta) \|u\|_{L^\infty(D)}, \\ \hat{R}(k)u(x) &= \int_D R(x, y, k) u(y) dy, \quad x \in \Omega, \\ k &\in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}), \end{aligned} \quad (4.7a)$$

$$\begin{aligned}
\|\hat{R}_\gamma(k)u\|_{C^{1+\delta}(\Omega)} &\leq c_3(D, \Omega, v, k, \gamma, \delta)\|u\|_{\mathbb{L}^\infty(D)}, \\
\hat{R}_\gamma(k)u(x) &= \int_D R_\gamma(x, y, k)u(y)dy, \quad x \in \Omega, \\
\gamma &\in \mathbb{S}^{d-1}, \quad k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma),
\end{aligned} \tag{4.7b}$$

for  $u \in \mathbb{L}^\infty(D)$ ,  $\delta \in [0, 1)$ , where  $\Omega$  is such an open bounded domain in  $\mathbb{R}^d$  that  $\bar{D} \in \Omega$  and  $C^{1+\delta}$  denotes  $C^1$  with the first derivatives belonging to the Hölder space  $C^\delta$ .

We will use also the Green formula:

$$\int_{\partial D} \left( \phi_1 \frac{\partial \phi_2}{\partial \nu} - \phi_2 \frac{\partial \phi_1}{\partial \nu} \right) dx = \int_D (\phi_1 \Delta \phi_2 - \phi_2 \Delta \phi_1) dx, \tag{4.8}$$

where  $\phi_1$  and  $\phi_2$  are arbitrary sufficiently regular functions in  $\bar{D}$ .

## 4.2 Proof of Theorem 3.1 and Proposition 3.1

For the case when  $\sin \alpha = 0$ , Theorem 3.1 and Proposition 3.1 were proved in [40]. In this subsection we generalize the proof of [40] to the case  $\sin \alpha \neq 0$ . We proceed from the following formulas and equations (being valid under assumption (2.2) on  $v^0$  and  $v$ ):

$$h(k, l) - h^0(k, l) = \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} \psi^0(x, -l)(v(x) - v^0(x))\psi(x, k)dx, \tag{4.9}$$

$$k, l \in \mathbb{C}^d \setminus (\mathcal{E}^0 \cup \mathcal{E}), \quad k^2 = l^2, \quad |\operatorname{Im} k| = |\operatorname{Im} l| \neq 0,$$

$$\psi(x, k) = \psi^0(x, k) + \int_{\mathbb{R}^d} R^0(x, y, k)(v(y) - v^0(y))\psi(y, k)dy, \tag{4.10}$$

$$x \in \mathbb{R}^d, \quad k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}^0),$$

where (4.10) at fixed  $k$  is considered as an equation for  $\psi = e^{ikx}\mu(x, k)$  with  $\mu \in \mathbb{L}^\infty(\mathbb{R}^d)$ ;

$$f(k, l) - f^0(k, l) = \left( \frac{1}{2\pi} \right)^d \int_{\mathbb{R}^d} \psi^{+,0}(x, -l)(v(x) - v^0(x))\psi^+(x, k)dx, \tag{4.11}$$

$$k, l \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^{+,0} \cup \mathcal{E}^+), \quad k^2 = l^2,$$

$$\psi^+(x, k) = \psi^{+,0}(x, k) + \int_{\mathbb{R}^d} R^{+,0}(x, y, k)(v(y) - v^0(y))\psi^+(y, k)dy, \quad (4.12)$$

$$x \in \mathbb{R}^d, \quad k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^{+,0}),$$

where (4.12) at fixed  $k$  is an equation for  $\psi^+ \in \mathbb{L}^\infty(\mathbb{R}^d)$ ;

$$h_\gamma(k, l) - h_\gamma^0(k, l) = \left(\frac{1}{2\pi}\right)^d \int_{\mathbb{R}^d} \psi_{-\gamma}^0(x, -k, -l)(v(x) - v^0(x))\psi_\gamma(x, k)dx, \\ \gamma \in \mathbb{S}^{d-1}, k \in \mathbb{R}^d \setminus (\mathcal{E}_\gamma^0 \cup \mathcal{E}_\gamma), \quad l \in \mathbb{R}^d, \quad k^2 = l^2, \quad (4.13)$$

$$\psi_\gamma(x, k) = \psi_\gamma^0(x, k) + \int_{\mathbb{R}^d} R_\gamma^0(x, y, k)(v(y) - v^0(y))\psi_\gamma(y, k)dy, \quad (4.14) \\ x \in \mathbb{R}^d, \quad \gamma \in \mathbb{S}^{d-1}, \quad k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma^0),$$

where (4.14) at fixed  $\gamma$  and  $k$  is considered as an equation for  $\psi_\gamma \in \mathbb{L}^\infty(\mathbb{R}^d)$ .

We recall that  $\psi^+$ ,  $f$ ,  $\psi$ ,  $h$ ,  $\psi_\gamma$ ,  $h_\gamma$  were defined in Sections 2, 3 by means of (2.3) - (2.9), (3.5). Equation (4.12) is well-known in the classical scattering theory for the Schrödinger equation (2.1). Formula (4.11) was given, in particular, in [50]. To our knowledge formula and equations (4.9), (4.10), (4.14) were given for the first time in [38], whereas formula (4.13) was given for the first time in [40].

In addition, under assumption (2.2) on  $v^0$  and  $v$ :

$$\text{equation (4.10) at fixed } k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}^0) \text{ is uniquely solvable} \\ \text{for } \psi = e^{ikx}\mu(x, k) \text{ with } \mu \in \mathbb{L}^\infty(\mathbb{R}^d) \text{ if and only if } k \notin \mathcal{E}; \quad (4.15a)$$

$$\text{equation (4.12) at fixed } k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}^{+,0}) \text{ is uniquely} \\ \text{solvable for } \psi^+ \in \mathbb{L}^\infty(\mathbb{R}^d) \text{ if and only if } k \notin \mathcal{E}^+; \quad (4.15b)$$

$$\text{equation (4.14) at fixed } \gamma \in \mathbb{S}^{d-1} \text{ and } k \in \mathbb{R}^d \setminus (\{0\} \cup \mathcal{E}_\gamma^+) \\ \text{is uniquely solvable for } \psi_\gamma \in \mathbb{L}^\infty(\mathbb{R}^d) \text{ if and only if } k \notin \mathcal{E}_\gamma. \quad (4.15c)$$

Let us prove Theorem 3.1 for the case of the Faddeev functions  $\psi$ ,  $h$ . The proof of Theorem 3.1 for the case of  $\psi^+$ ,  $f$  and the proof of Proposition 3.1 are similar.

Note that formula (3.1) follows directly from (4.6) and (4.9).



Using (2.17) and applying (4.6) for equation (4.10), we get that

$$\psi(x, k) - \psi^0(x, k) = \int_{\partial D} \int_{\partial D} [R^0(x, \xi, k)]_{\xi, \alpha} (M_{\alpha, v} - M_{\alpha, v^0})(\xi, y, E) [\psi(y, k)]_{\alpha} d\xi dy, \quad x \in \mathbb{R}^d \setminus \bar{D}, \quad (4.16)$$

where

$$[R^0(x, \xi, k)]_{\xi, \alpha} = \left( \cos \alpha - \sin \alpha \frac{\partial}{\partial \nu_{\xi}} \right) R^0(x, \xi, k). \quad (4.17)$$

Equation (3.2) follows from formula (4.16), definition (1.5) and the property that

$$\lim_{\varepsilon \rightarrow +0} \left( \cos \alpha - \sin \alpha \frac{\partial}{\partial \nu_x} \right) u(x + \varepsilon \nu_x) = [u(x)]_{\alpha}, \quad x \in \partial D, \quad (4.18)$$

for  $u(x) = \psi(x, k) - \psi^0(x, k)$ .

### 4.3 Proofs of Propositions 3.2 and 3.4

In this subsection we prove Propositions 3.2, 3.4 for the case of equation (3.2) for  $[\psi]_{\alpha}$ . The proofs of Propositions 3.2 and 3.4 for the cases of  $\psi^+$  and  $\psi_{\gamma}$  are absolutely similar.

*Proof of Proposition 3.2.* The proof of Proposition 3.2 for the case of  $\sin \alpha = 0$  was given in [40]. Let us assume that  $\sin \alpha \neq 0$ .

Using (4.2), we find that

$$(M_{\alpha, v} - M_{\alpha, v^0})(\xi, y, E) = \frac{1}{\sin^2 \alpha} (G_{\alpha, v} - G_{\alpha, v^0})(\xi, y, E), \quad \xi, y \in \partial D. \quad (4.19)$$

Using (2.17), (4.1), (4.8) and the impedance boundary condition (1.6) for

$G_{\alpha,v}$ ,  $G_{\alpha,v^0}$ , we get that

$$\begin{aligned}
& \int_{\partial D} [R^0(x, \xi, k)]_{\alpha, \xi} (G_{\alpha,v} - G_{\alpha,v^0})(\xi, y, E) d\xi = \\
& = \int_{\partial D} \left( [R^0(x, \xi, k)]_{\alpha, \xi} (G_{\alpha,v} - G_{\alpha,v^0})(\xi, y, E) d\xi - \right. \\
& \quad \left. - R^0(x, \xi, k) [(G_{\alpha,v} - G_{\alpha,v^0})(\xi, y, E)]_{\alpha, \xi} \right) d\xi = \\
& = \sin \alpha \int_D \left( R^0(x, \xi, k) \Delta_\xi (G_{\alpha,v} - G_{\alpha,v^0})(\xi, y, E) d\xi - \right. \\
& \quad \left. - (G_{\alpha,v} - G_{\alpha,v^0})(\xi, y, E) \Delta_\xi R^0(x, \xi, k) \right) d\xi = \\
& = \sin \alpha \int_D R^0(x, \xi, k) (v(\xi) - v^0(\xi)) G_{\alpha,v}(\xi, y, E) d\xi, \quad x \in \mathbb{R}^d \setminus \bar{D}, \quad y \in \partial D.
\end{aligned} \tag{4.20}$$

Combining (4.16), (4.19) and (4.20), we obtain that

$$A_\alpha(x, y, k) = \lim_{\varepsilon \rightarrow +0} \left( \cos \alpha - \sin \alpha \frac{\partial}{\partial \nu_x} \right) B_\alpha(x + \varepsilon \nu_x, y, k), \quad x, y \in \partial D, \tag{4.21}$$

where

$$\begin{aligned}
B_\alpha(x, y, k) &= \int_{\partial D} [R^0(x, \xi, k)]_{\alpha, \xi} (M_{\alpha,v} - M_{\alpha,v^0})(\xi, y, E) d\xi = \\
&= \frac{1}{\sin \alpha} \int_D R^0(x, \xi, k) (v(\xi) - v^0(\xi)) G_{\alpha,v}(\xi, y, E) d\xi, \tag{4.22} \\
&\quad x \in \mathbb{R}^d \setminus \bar{D}, \quad y \in \partial D.
\end{aligned}$$

Thus, we have that the limit in (4.21) (and, hence, in (3.3)) is well defined and

$$A_\alpha(x, y, k) = \frac{1}{\sin \alpha} \int_D [R^0(x, \xi, k)]_{x, \alpha} (v(\xi) - v^0(\xi)) G_{\alpha,v}(\xi, y, E) d\xi, \quad x, y \in \partial D. \tag{4.23}$$

Let  $\hat{A}_\alpha(k)$  denote the linear integral operator on  $\partial D$  with the Schwartz kernel  $A_\alpha(x, y, k)$  of (3.3), (4.23). Using (4.5), (4.7), (4.23), we obtain that

$$\begin{aligned} \hat{A}_\alpha(k) : \mathbb{L}^\infty(\partial D) &\rightarrow C^\delta(\partial D) \\ &\text{is a bounded linear operator.} \end{aligned} \tag{4.24}$$

As a corollary of (4.24),  $\hat{A}_\alpha(k)$  is a compact operator in  $\mathbb{L}^\infty(D)$ . ■

*Proof of Proposition 3.4.* Using (2.17), (3.7) and (4.8), we get that

$$\begin{aligned} &\int_{\partial D} \left( \phi_\alpha(\xi, y) \frac{\partial}{\partial \nu_\xi} R^0(x, \xi, k) - R^0(x, \xi, k) \frac{\partial}{\partial \nu_\xi} \phi_\alpha(\xi, y) \right) d\xi = \\ &= \int_D (\phi_\alpha(\xi, y) \Delta_\xi R^0(x, \xi, k) - R^0(x, \xi, k) \Delta_\xi \phi_\alpha(\xi, y)) d\xi = \\ &= \int_D R^0(x, \xi, k) (v^0(\xi) - E + \lambda) \phi_\alpha(\xi, y) d\xi, \\ &\quad x \in \mathbb{R}^d \setminus \bar{D}, \ y \in \partial D. \end{aligned} \tag{4.25}$$

Combining (3.7), (4.22) and (4.25), we find that

$$\begin{aligned} B_\alpha(x, y, k) &= \int_{\partial D} [R^0(x, \xi, k)]_{\xi, \alpha} \phi_\alpha(\xi, y) d\xi = \\ &= \int_{\partial D} R^0(x, \xi, k) [\phi_\alpha(\xi, y)]_{\xi, \alpha} d\xi - \sin \alpha \int_D R^0(x, \xi, k) (v^0(\xi) - E + \lambda) \phi_\alpha(\xi, y) d\xi, \\ &\quad x \in \mathbb{R}^d \setminus \bar{D}, \ y \in \partial D. \end{aligned} \tag{4.26}$$

Combining (4.21) and (4.26), we obtain (3.8).

Formula (3.10) follows from (3.7) and the definition of  $\hat{\Phi}$ . ■

## 5 Proof of Proposition 3.3

For the case when  $\sin \alpha = 0$ , Proposition 3.3 was proved in [40]. In this section we prove Proposition 3.3 for  $\sin \alpha \neq 0$ . We will prove Proposition 3.3

for the case of equation (3.2) for  $[\psi]_\alpha$ . The proofs for the cases of  $\psi^+$  and  $\psi_\gamma$  are similar.

According to (4.15), to prove Proposition 3.3 (for the case of  $\psi$ ) it is sufficient to show that equation (3.2) (at fixed  $k \in \mathbb{C}^d \setminus (\mathbb{R}^d \cup \mathcal{E}^0)$ ) is uniquely solvable in the space of bounded functions on  $\partial D$  if and only if equation (4.10) is uniquely solvable for  $\psi = e^{ikx}\mu(x, k)$  with  $\mu \in \mathbb{L}^\infty(\mathbb{R}^d)$ .

Let equation (4.10) have several solutions. Then, repeating the proof of Theorem 3.1 separately for each solution, we find that  $[\psi]_\alpha$  on  $\partial D$  for each of these solutions satisfies equation (3.2). Thus, using also (1.7) we obtain that equation (3.2) has at least as many solutions as equation (4.10).

To prove the converse (and thereby to prove Proposition 3.3) it remains to show that any solution  $[\psi]_\alpha$  of (3.2) can be continued to a continuous solution of (4.10).

Let  $\psi$  be the solution of (1.1) with the impedance boundary data  $[\psi]_\alpha$ , satisfying (3.2). Let

$$\psi_1(x) = \psi^0(x, k) + \int_D R^0(x, y, k)(v(y) - v^0(y))\psi(y)dy, \quad x \in \mathbb{R}^d. \quad (5.1)$$

Using (4.7), we obtain that

$$\psi_1 \text{ defined by (5.1) belongs to } C^{1+\delta}(\mathbb{R}^d), \delta \in [0, 1). \quad (5.2)$$

We have that

$$(-\Delta + v^0(x) - E)\psi(x) = (v^0(x) - v(x))\psi(x), \quad x \in D, \quad (5.3)$$

$$\begin{aligned} (-\Delta + v^0(x) - E)\psi_1(x) &= \int_D -\delta(x - y)(v(y) - v^0(y))\psi(y)dy = \\ &= (v^0(x) - v(x))\psi(x), \quad x \in D. \end{aligned} \quad (5.4)$$

Combining (4.6) and (4.22), we get that

$$\int_D R^0(x, y, k)(v(y) - v^0(y))\psi(y)dy = \int_{\partial D} B_\alpha(x, y, k)[\psi(y)]_\alpha dy, \quad x \in \mathbb{R}^d \setminus \bar{D}. \quad (5.5)$$

Using (3.2), (4.21), (5.2), (5.5), we find that

$$[\psi_1(x)]_\alpha = [\psi^0(x, k)]_\alpha + \int_{\partial D} A_\alpha(x, y, k)[\psi(y)]_\alpha dy = [\psi(x)]_\alpha, \quad x \in \partial D. \quad (5.6)$$

Using (5.3), (5.4) and (5.6), we obtain that

$$\begin{aligned} (-\Delta + v^0(x) - E)(\psi_1(x) - \psi(x)) &= 0, \quad x \in D, \\ [\psi_1(x) - \psi(x)]_\alpha &= 0, \quad x \in \partial D. \end{aligned} \quad (5.7)$$

Since  $v^0$  satisfies (1.7), we get that

$$\psi_1(x) = \psi(x), \quad x \in \bar{D}. \quad (5.8)$$

Combining (5.1), (5.2) and (5.8), we find that  $\psi_1$  is a continuous solution of (4.10).

## 6 Proofs of properties (4.4), (4.5)

As it was mentioned in Subsection 4.1, properties (4.4), (4.5) are well-known for  $\cos \alpha = 0$  (the case of the Neumann boundary condition). To extend these properties to the case of general  $\alpha$ ,  $v$ ,  $E$ , we use the following schema:

1.  $G_{\alpha_1, v} \rightarrow G_{\alpha_2, v}$  by means of Lemma 6.1 given below (with  $\sin \alpha_1 \neq 0$  and  $\sin \alpha_2 \neq 0$ ),
2.  $G_{\alpha, v_1} \rightarrow G_{\alpha, v_2}$  by means of Lemma 6.2 given below.

The proofs of steps 1, 2 are based on the theory of Fredholm linear integral equations of the second kind.

Starting from (4.4), (4.5) for  $\cos \alpha = 0$  and combining steps 1, 2 and the property

$$G_{\alpha, v}(\cdot, \cdot, E) = G_{\alpha, v-E}(\cdot, \cdot, 0), \quad (6.1)$$

we obtain these properties for the case when  $\sin \alpha \neq 0$ .

As it was already mentioned in Section 4, properties (4.4), (4.5) are well-known for  $\sin \alpha = 0$  (the case of the Dirichlet boundary condition).

**Lemma 6.1.** *Let  $D$  satisfy (1.2) and potential  $v$  satisfy (1.3), (1.7) for some fixed  $E$  and for  $\alpha = \alpha_1$ ,  $\alpha = \alpha_2$  simultaneously, where  $\sin \alpha_1 \neq 0$  and  $\sin \alpha_2 \neq 0$ . Let  $G_j$  denote the Green function  $G_{\alpha_j, v}$ ,  $j = 1, 2$ . Let  $G_1$  satisfy:*

$$G_1(x, y, E) \text{ is continuous in } x, y \in \bar{D}, \quad x \neq y, \quad (6.2)$$

$$\begin{aligned} |G_1(x, y, E)| &\leq a_1 |x - y|^{2-d} \quad \text{for } d \geq 3, \\ |G_1(x, y, E)| &\leq a_1 |\ln |x - y|| \quad \text{for } d = 2, \\ &x, y \in \bar{D}. \end{aligned} \quad (6.3)$$

Then:

$$G_2(x, y, E) \text{ is continuous in } x, y \in \bar{D}, x \neq y, \quad (6.4)$$

$$\begin{aligned} |G_2(x, y, E)| &\leq a_2 |x - y|^{2-d} \quad \text{for } d \geq 3, \\ |G_2(x, y, E)| &\leq a_2 |\ln |x - y|| \quad \text{for } d = 2, \\ &x, y \in \bar{D}, \end{aligned} \quad (6.5)$$

where  $a_2 = a_2(D, E, a_1, v, \alpha_1, \alpha_2) > 0$ .

*Proof of Lemma 6.1.* First, we derive formally some formulas and equations relating the Green functions  $G_1$  and  $G_2$ . Then, proceeding from these formulas and equations, we obtain, in particular, estimates (6.4), (6.5).

Consider  $W = G_2 - G_1$ . Using definitions of  $G_1$ ,  $G_2$  and formula (4.3), we find that:

$$(-\Delta_x + v(x) - E)W(x, y) = 0, \quad x, y \in D, \quad (6.6)$$

$$\begin{aligned} \left( \cos \alpha_2 W(x, y) - \sin \alpha_2 \frac{\partial W}{\partial \nu_x}(x, y) \right) \Big|_{x \in \partial D} &= \\ &= - \left( \cos \alpha_2 G_1(x, y, E) - \sin \alpha_2 \frac{\partial G_1}{\partial \nu_x}(x, y, E) \right) \Big|_{x \in \partial D} = \\ &= - \left( \cos \alpha_2 G_1(x, y, E) - \sin \alpha_2 \frac{\cos \alpha_1}{\sin \alpha_1} G_1(x, y, E) \right) \Big|_{x \in \partial D} = \\ &= \frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_1} G_1(x, y, E) \Big|_{x \in \partial D}, \quad y \in D, \end{aligned} \quad (6.7)$$

$$W(x, y) = \frac{1}{\sin \alpha_1} \int_{\partial D} \left( \cos \alpha_1 W(\xi, y) - \sin \alpha_1 \frac{\partial W}{\partial \nu_\xi}(\xi, y) \right) G_1(\xi, x, E) d\xi, \quad x, y \in D. \quad (6.8)$$

Using (6.7) and (6.8), we find the following linear integral equation for  $W(\cdot, y)$  on  $\partial D$ :

$$W(\cdot, y) = W_0(\cdot, y) + \hat{K}_1 W(\cdot, y), \quad y \in D, \quad (6.9)$$

where

$$W_0(x, y) = \frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_2} \int_{\partial D} G_1(\xi, x, E) G_1(\xi, y, E) d\xi, \quad (6.10)$$

$$\hat{K}_1 u(x) = \frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_2 \sin \alpha_1} \int_{\partial D} G_1(\xi, x, E) u(\xi) d\xi, \quad (6.11)$$

$x \in \partial D, y \in D, u$  is a test function.

In addition, for

$$\delta_n W = W - \sum_{j=1}^n (\hat{K}_1)^{j-1} W_0 \quad (6.12)$$

equation (6.9) takes the form

$$\delta_n W = (\hat{K}_1)^n W_0 + \hat{K}_1 \delta_n W. \quad (6.13)$$

Our analysis based on (6.6)-(6.13) is given bellow.

Using (6.2), (6.3), we obtain that

$$(\hat{K}_1)^n W_0 \in C(\partial D \times \bar{D}) \text{ for sufficiently great } n \text{ with respect to } d, \quad (6.14)$$

$$\hat{K}_1 \text{ is a compact operator in } C(\partial D). \quad (6.15)$$

Let us show that the homogeneous equation

$$u = \hat{K}_1 u, \quad u \in C(\partial D), \quad (6.16)$$

has only trivial solution  $u \equiv 0$ .

Using the fact that the potential  $v$  satisfy (1.7) for  $\alpha = \alpha_1$ , we define  $\psi$  by

$$\begin{aligned} (-\Delta + v(x) - E)\psi(x) &= 0, \quad x \in D, \\ \cos \alpha_1 \psi|_{\partial D} - \sin \alpha_1 \frac{\partial \psi}{\partial \nu}|_{\partial D} &= u. \end{aligned} \quad (6.17)$$

Due to (4.3), we have that

$$\psi(x) = \frac{1}{\sin \alpha_1} \int_{\partial D} (\cos \alpha_1 \psi(\xi) - \sin \alpha_1 \frac{\partial \psi}{\partial \nu}(\xi)) G_1(\xi, x, E) d\xi, \quad x \in D. \quad (6.18)$$

Using (6.16), (6.18), we find that

$$\frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_2} \psi(x) = \hat{K}_1 u(x) = u(x), \quad x \in \partial D. \quad (6.19)$$

Therefore, we have that

$$\cos \alpha_1 \psi(x) - \sin \alpha_1 \frac{\partial \psi}{\partial \nu}(x) = \frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_2} \psi(x), \quad x \in \partial D. \quad (6.20)$$

Since  $\sin \alpha_1 \neq 0$  and  $\sin \alpha_2 \neq 0$ , using (6.20), we obtain that

$$\cos \alpha_2 \psi(x) - \sin \alpha_2 \frac{\partial \psi}{\partial \nu}(x) = 0 \quad (6.21)$$

Taking into account the fact that the potential  $v$  satisfy (1.7) for  $\alpha = \alpha_2$ , we get that  $\psi \equiv 0$  and  $u \equiv 0$ .

Proceeding from

$$F = W(x, y) \quad \text{and} \quad F' = \frac{\cos \alpha_2}{\sin \alpha_2} W(x, y) - \frac{\sin(\alpha_2 - \alpha_1)}{\sin \alpha_1 \sin \alpha_2} G_1(x, y, E), \quad (6.22)$$

$$x \in \partial D, \quad y \in \bar{D},$$

found from (6.9), (6.13) and (6.7) (with  $F'$  substituted in place of  $\partial W / \partial \nu_x$ ), we consider

$$W(x, y) = \frac{1}{\sin \alpha_1} \int_{\partial D} \left( \cos \alpha_1 F(\xi, y) - \sin \alpha_1 F'(\xi, y) \right) G_1(\xi, x, E) d\xi, \quad x, y \in \bar{D}. \quad (6.23)$$

Using (6.9) and properties of  $G_1$  (including formula (4.3)), we subsequently obtain that

$$\lim_{\varepsilon \rightarrow +0} W(x - \varepsilon \nu_x, y) = F(x, y), \quad x \in \partial D, \quad y \in \bar{D}, \quad (6.24)$$

$$W \text{ satisfies (6.6),} \quad (6.25)$$

$$\lim_{\varepsilon \rightarrow +0} \frac{\partial}{\partial \nu_x} W(x - \varepsilon \nu_x, y) = F'(x, y), \quad x \in \partial D, \quad y \in \bar{D}. \quad (6.26)$$

From (6.2), (6.3), (6.10)-(6.16), (6.24)-(6.26) it follows that  $G_2$  defined as  $G_2 = G_1 + W$  is the Green function for the operator  $\Delta - v + E$  in  $D$  with the impedance boundary condition (1.6) for  $\alpha = \alpha_2$  and that  $G_2$  satisfies (6.4), (6.5). ■

**Lemma 6.2.** *Let  $D$  satisfy (1.2) and potentials  $v_1, v_2$  satisfy (1.3), (1.7) for some fixed  $E$  and  $\alpha$ . Let  $G_j$  denote the Green function  $G_{\alpha, v_j}$ ,  $j = 1, 2$ . Let  $G_1$  satisfy:*

$$G_1(x, y, E) \text{ is continuous in } x, y \in \bar{D}, \quad x \neq y, \quad (6.27)$$



$$\begin{aligned}
|G_1(x, y, E)| &\leq a_3 |x - y|^{2-d} \quad \text{for } d \geq 3, \\
|G_1(x, y, E)| &\leq a_3 |\ln |x - y|| \quad \text{for } d = 2, \\
x, y &\in \bar{D}.
\end{aligned} \tag{6.28}$$

Then:

$$G_2(x, y, E) \text{ is continuous in } x, y \in \bar{D}, x \neq y, \tag{6.29}$$

$$\begin{aligned}
|G_2(x, y, E)| &\leq a_4 |x - y|^{2-d} \quad \text{for } d \geq 3, \\
|G_2(x, y, E)| &\leq a_4 |\ln |x - y|| \quad \text{for } d = 2, \\
x, y &\in \bar{D},
\end{aligned} \tag{6.30}$$

where  $a_4 = a_4(D, E, a_3, v_1, v_2, \alpha) > 0$ .

*Proof of Lemma 6.2.* First, we derive formally some formulas and equations relating the Green functions  $G_1$  and  $G_2$ . Then, proceeding from these formulas and equations, we obtain, in particular, estimates (6.29), (6.30).

Using (4.1), the impedance boundary condition for  $G_1, G_2$ , we find that

$$\begin{aligned}
G_1(x, y, E) &= \int_D G_1(x, \xi, E) (\Delta_\xi - v_2(\xi) + E) G_2(\xi, y, E) d\xi, \\
G_2(x, y, E) &= \int_D G_2(\xi, y, E) (\Delta_\xi - v_1(\xi) + E) G_1(x, \xi, E) d\xi, \\
\int_{\partial D} \left( G_1(x, \xi, E) \frac{\partial G_2}{\partial \nu_\xi}(\xi, y, E) - G_2(\xi, y, E) \frac{\partial G_1}{\partial \nu_\xi}(x, \xi, E) \right) d\xi &= 0, \\
x, y &\in D.
\end{aligned} \tag{6.31}$$

Combining (6.31) with (4.8), we get that

$$G_2(\cdot, y, E) - G_1(\cdot, y, E) = \hat{K}_2 G_2(\cdot, y, E), \quad y \in D, \tag{6.32}$$

where

$$\hat{K}_2 u(x) = \int_D (v_2(\xi) - v_1(\xi)) G_1(x, \xi, E) u(\xi) d\xi. \tag{6.33}$$

In addition, for

$$\delta_n G = G_2 - \sum_{j=1}^n (\hat{K}_2)^{j-1} G_1 \tag{6.34}$$

equation (6.32) takes the form

$$\delta_n G = (\hat{K}_2)^n G_1 + \hat{K}_2 \delta_n G. \quad (6.35)$$

Our analysis based on (6.31)-(6.35) is given bellow.

Using (6.27), (6.28), we find that

$$(\hat{K}_2)^n G_1 \in C(\bar{D} \times \bar{D}) \text{ for sufficiently great } n \text{ with respect to } d, \quad (6.36)$$

$$\hat{K}_2 \text{ is a compact operator in } C(\bar{D}). \quad (6.37)$$

Let us show that the homogeneous equation

$$u = \hat{K}_2 u, \quad u \in C(\bar{D}), \quad (6.38)$$

has only trivial solution  $u \equiv 0$ . Using (6.33), (6.38) and properties of the Green function  $G_1$ , we find that

$$\begin{aligned} (-\Delta + v_1(x) - E)u(x) &= \int_D -\delta(x - \xi) (v_2(\xi) - v_1(\xi)) u(\xi) d\xi = \\ &= (v_1 - v_2)u(x), \quad x \in D, \\ \cos \alpha u(x) - \sin \alpha \frac{\partial u}{\partial \nu}(x) &= 0, \quad x \in \partial D. \end{aligned} \quad (6.39)$$

Using (6.27), (6.28), we find that  $u \in C(\bar{D})$ . Taking into account the fact that the potential  $v_2$  satisfy (1.7), we get that  $u \equiv 0$ .

Proceeding from (6.27), (6.28), (6.36), (6.37) it follows that  $G_2$  found from (6.32), (6.35) is the Green function for the operator  $\Delta - v + E$  in  $D$  with the impedance boundary condition (1.6) for  $v = v_2$  and that  $G_2$  satisfies (6.29), (6.30). ■

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## References

- [1] G.Alessandrini, *Stable determination of conductivity by boundary measurements*, Appl. Anal. 27, 1988, 153-172.
- [2] G. Alessandrini, S. Vassella, *Lipschitz stability for the inverse conductivity problem*, Adv. in Appl. Math. 35, 2005, no.2, 207-241.
- [3] N.V. Alexeenko, V.A. Burov and O.D. Rumyantseva, *Solution of the three-dimensional acoustical inverse scattering problem. The modified Novikov algorithm*, Acoust. J. 54(3), 2008, 469-482 (in Russian), English transl.: Acoust. Phys. 54(3), 2008, 407-419.
- [4] L. Beilina, M.V. Klibanov, *Approximate global convergence and adaptivity for coefficient inverse problems*, Springer (New York), 2012. 407 pp.
- [5] H. Begehr and T. Vaitekhovich, *Some harmonic Robin functions in the complex plane*, Adv. Pure Appl. Math. 1, 2010, 19-34.
- [6] F. A. Berezin and M. A. Shubin, *The Schrödinger Equation*, Vol. 66 of Mathematics and Its Applications, Kluwer Academic, Dordrecht, 1991.
- [7] Yu.M. Berezanskii, *The uniqueness theorem in the inverse problem of spectral analysis for the Schrodinger equation*. (Russian) Trudy Moskov. Mat. Obsc. 7 (1958) 1-62.
- [8] J. Bikowski, K. Knudsen, J.L. Mueller, *Direct numerical reconstruction of conductivities in three dimensions using scattering transforms*. Inverse Problems 27(1), 2011, 015002, 19 pp.
- [9] A.V. Bogatyrev, V.A. Burov, S.A. Morozov, O.D. Rumyantseva, E.G. Sukhov, *Numerical realization of algorithm for exact solution of two-dimensional monochromatic inverse problem of acoustical scattering*, Acoust. Imaging 25, 2000, 65-70.
- [10] A. L. Bukhgeim, *Recovering a potential from Cauchy data in the two-dimensional case*, J. Inverse Ill-Posed Probl. 16(1), 2008, 19-33.
- [11] V.A. Burov, N.V. Alekseenko, O.D. Rumyantseva, *Multifrequency generalization of the Novikov algorithm for the two-dimensional inverse scattering problem*, Acoustical Physics 55, 2009, no. 6, 843-856.

- [12] V.A. Burov, A.S. Shurup, O.D. Rumyantseva, D.I. Zotov, *Functional-analytical solution of the problem of acoustic tomography from point transducer data*, Izvestiya Rossiiskoi Akademii Nauk. Seriya Fizicheskaya 76(12),2012, 1524-1529 (in Russian)
- [13] A.P. Calderón, *On an inverse boundary problem*, Seminar on Numerical Analysis and its Applications to Continuum Physics, Soc. Brasileira de Matematica, Rio de Janeiro, 1980, 61-73.
- [14] G. Eskin, *Lectures on Linear Partial Differential Equations*, Graduate Studies in Mathematics, Vol.123, American Mathematical Society, 2011
- [15] V.L. Druskin, *The unique solution of the inverse problem in electrical surveying and electrical well logging for piecewise-constant conductivity*, Physics of the Solid Earth 18(1), 1982, 51-53.
- [16] L.D. Faddeev, *Growing solutions of the Schrödinger equation*, Dokl. Akad. Nauk SSSR, 165, N.3, 1965, 514-517 (in Russian); English Transl.: Sov. Phys. Dokl. 10, 1966, 1033-1035.
- [17] L.D. Faddeev, *The inverse problem in the quantum theory of scattering. II*, Current problems in mathematics, Vol. 3, 1974, pp. 93-180, 259. Akad. Nauk SSSR Vsesojuz. Inst. Nauch. i Tehn. Informacii, Moscow(in Russian); English Transl.: J.Sov. Math. 5, 1976, 334-396.
- [18] I.M. Gelfand, *Some problems of functional analysis and algebra*, Proceedings of the International Congress of Mathematicians, Amsterdam, 1954, pp.253-276.
- [19] F. Gesztesy and M. Mitrea, *Robin-to-Robin Maps and Krein-Type Resolvent Formulas for Schrodinger Operators on Bounded Lipschitz Domains*, Modern Analysis and Applications, Operator Theory: Advances and Applications, Volume 191, 2009, Part 1, 81-113.
- [20] P.G. Grinevich, *The scattering transform for the two-dimensional Schrödinger operator with a potential that decreases at infinity at fixed nonzero energy*, Uspekhi Mat. Nauk 55:6(336),2000, 3-70 (Russian); English translation: Russian Math. Surveys 55:6, 2000, 1015-1083.

- [21] P.G. Grinevich, S.P. Novikov, *Two-dimensional "inverse scattering problem" for negative energies and generalized-analytic functions. 1. Energies below the ground state*, Funkt. Anal. Prilozhen. 22:1, 1988, 23-33 (Russian); English translation: Funct. Anal. Appl. 22, 1988, 19-27.
- [22] P.G. Grinevich, R.G. Novikov, *Faddeev eigenfunctions for multipoint potentials*, e-print: arXiv: 1211.0292.
- [23] G.M. Henkin and R.G. Novikov, *The  $\bar{\partial}$ -equation in the multidimensional inverse scattering problem*, Uspekhi Mat. Nauk 42(3), 1987, 93-152 (in Russian); English Transl.: Russ. Math. Surv. 42(3), 1987, 109-180.
- [24] M.I. Isaev, *Exponential instability in the Gel'fand inverse problem on the energy intervals*, J. Inverse Ill-Posed Probl., 19(3), 2011, 453-473.
- [25] M.I. Isaev, *Instability in the Gel'fand inverse problem at high energies*, Applicable Analysis (to appear), e-print arXiv:1206.2328.
- [26] M.I. Isaev, R.G. Novikov *Stability estimates for determination of potential from the impedance boundary map*, Algebra and Analysis, 25, 2013, e-print arXiv:1112.3728v2.
- [27] M.I. Isaev, R.G. Novikov *Energy and regularity dependent stability estimates for the Gel'fand inverse problem in multidimensions*, J. of Inverse and Ill-posed Probl., 20(3), 2012, 313-325.
- [28] R. Kohn, M. Vogelius, *Determining conductivity by boundary measurements II*, Interior results, Comm. Pure Appl. Math. 38, 1985, 643-667.
- [29] L. Lanzani and Z. Shen, *On the Robin boundary condition for Laplace's equation in Lipschitz domains*, Comm. Partial Differential Equations, 29, 2004, 91-109.
- [30] R.B. Lavine and A.I. Nachman, *On the inverse scattering transform of the  $n$ -dimensional Schrödinger operator* Topics in Soliton Theory and Exactly Solvable Nonlinear Equations ed M Ablowitz, B Fuchssteiner and M Kruskal (Singapore: World Scientific), 1987, pp 33-44.
- [31] M.M. Lavrent'ev, V.G. Romanov, S.P. Shishatskii, *Ill-posed problems of mathematical physics and analysis*, Translated from the Russian by J. R. Schulenberger. Translation edited by Lev J. Leifman. Translations of

Mathematical Monographs, 64. American Mathematical Society, Providence, RI, 1986. vi+290 pp.

- [32] N. Mandache, *Exponential instability in an inverse problem for the Schrödinger equation*, Inverse Problems. 17, 2001, 1435-1444.
- [33] R.G. Newton, *Inverse Schrödinger scattering in three dimensions*. Texts and Monographs in Physics. Springer-Verlag, Berlin, 1989. x+170 pp.
- [34] A. Nachman, *Reconstructions from boundary measurements*, Ann. Math. 128, 1988, 531-576.
- [35] A. Nachman, *Global uniqueness for a two-dimensional inverse boundary value problem*, Ann. Math. 143, 1996, 71-96.
- [36] R.G. Novikov, *Multidimensional inverse spectral problem for the equation  $-\Delta\psi + (v(x) - Eu(x))\psi = 0$* , Funkt. Anal. Prilozhen. 22(4), 1988, 11-22 (in Russian); Engl. Transl.: Funct. Anal. Appl. 22, 1988, 263-72.
- [37] R.G. Novikov, *The inverse scattering problem on a fixed energy level for the two-dimensional Schrödinger operator*, J.Funct. Anal. 103 (1992), 409-463.
- [38] R.G. Novikov,  *$\bar{\partial}$ -method with nonzero background potential. Application to inverse scattering for the two-dimensional acoustic equation*, Comm. Partial Differential Equations 21, no. 3-4, 1996, 597-618.
- [39] R.G. Novikov, *Approximate solution of the inverse problem of quantum scattering theory with fixed energy in dimension 2*, Proceedings of the Steklov Mathematical Institute 225, 1999, Solitony Geom. Topol. na Perekrest., 301-318 (in Russian); Engl. Transl. in Proc. Steklov Inst. Math. 225, no. 2, 1999, 285-302.
- [40] R.G. Novikov, *Formulae and equations for finding scattering data from the Dirichlet-to-Neumann map with nonzero background potential*, Inverse Problems 21, 2005, 257-270.
- [41] R.G. Novikov, *The  $\bar{\partial}$ -approach to approximate inverse scattering at fixed energy in three dimensions*. IMRP Int. Math. Res. Pap. 2005, no. 6, 287-349.

- [42] R.G. Novikov, *On non-overdetermined inverse scattering at zero energy in three dimensions*, Ann. Scuola Norm. Sup. Pisa Cl. Sci. 5, 2006, 279-328.
- [43] R.G. Novikov, *An effectivization of the global reconstruction in the Gel'fand-Calderon inverse problem in three dimensions*, Contemporary Mathematics, 494, 2009, 161-184.
- [44] R.G. Novikov, *New global stability estimates for the Gel'fand-Calderon inverse problem*, Inverse Problems 27, 2011, 015001(21pp); e-print arXiv:1002.0153.
- [45] R.G. Novikov and M. Santacesaria, *A global stability estimate for the Gel'fand-Calderon inverse problem in two dimensions*, J.Inverse Ill-Posed Probl., 18(7), 2010, 765-785.
- [46] R.G. Novikov and M. Santacesaria, *Global uniqueness and reconstruction for the multi-channel Gel'fand-Calderon inverse problem in two dimensions*, Bulletin des Sciences Mathematiques 135(5), 2011, 421-434.
- [47] R.G. Novikov and M. Santacesaria, *Monochromatic Reconstruction Algorithms for Two-dimensional Multi-channel Inverse Problems*, International Mathematics Research Notes, 2012, doi: 10.1093/imrn/rns025.
- [48] L. Rondi, *A remark on a paper by Alessandrini and Vessella*, Adv. in Appl. Math. 36 (1), 2006, 67-69.
- [49] M. Santacesaria, *Global stability for the multi-channel Gel'fand-Calderon inverse problem in two dimensions*, Bulletin des Sciences Mathematiques, 136(7), 2012, 731-744.
- [50] P. Stefanov, *A uniqueness result for the inverse back-scattering problem*, Inverse Problems, 6, 1990, 1055-1064.
- [51] J. Sylvester and G. Uhlmann, *A global uniqueness theorem for an inverse boundary value problem*, Ann. of Math. 125, 1987, 153-169.
- [52] R. Weder, *Generalized limiting absorption method and multidimensional inverse scattering theory*, Mathematical Methods in the Applied Sciences, 14, 1991, 509-524.

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